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ARTICLE

# Remarkable Magnetic Properties of Co-doped ZnO Nanorods Array at Room Temperature Originated from Vertical Growth on Zn Foil

Li Yingying<sup>1</sup>, Dong Xiang<sup>2</sup>, Zhang Haiqian<sup>3</sup>

<sup>1</sup> Staff Room of Emergency Rescue, Public Security Fire Force College, Kunming 650208, China; <sup>2</sup> National Engineering Researching Centre of Waste Resources, Kunming University of Science and Technology, Kunming 650093, China; <sup>3</sup> Nanjing University of Aeronautics and Astronautics, Nanjing 210016, China

**Abstract:** Large-area arrays of highly oriented Co-doped ZnO nanorods with hexagonal structure were grown on Zn substrates by single-step hydrothermal process. Structure analysis indicates that the as-prepared products have wurtzite structure, and no other secondary phase is found in the nanorods. The intensity of UV emission peak decreases with the increase of Co doping concentration. When Co is doped into ZnO lattice, oxygen vacancies and Zn interstitials are created. The concentration of these defects increases with rise of the Co concentration. Magnetic measurement reveals that the Co doped ZnO nanorod arrays exhibit obvious room-temperature ferromagnetic behavior with a large coercive field  $H_c$  of 52.8 kA/m, which makes them potentially useful as building components for spintronics.

Key words: semiconductor; magnetic material; nanorod array; ZnO; hydrothermal

Diluted magnetic semiconductors (DMSs) have attracted considerable research interest in recent years due to their great potential applications in spintronic devices, such as spin field-effect transistors, non-volatile memory devices and quantum computer <sup>[1,2]</sup>. A key work to realize spintronic devices is to develop ferromagnetic DMSs with a Curie temperature above room temperature. Transition metal-doped ZnO DMSs have been extensively investigated since theoretic studies of Dietl et al. predicted their Curie temperature  $(T_c)$  to be above room temperature in 2000<sup>[3]</sup>. Following the original research by Ueda et al.<sup>[4]</sup>, many reports have been given about Co doped ZnO thin films exhibiting ferromagnetism with Tc above room temperature. Wu et al. <sup>[5]</sup> prepared Zn<sub>1-x</sub>Co<sub>x</sub>O nanowires using chemical vapor deposition (CVD) methods. Chen et al. <sup>[6]</sup> reported that room temperature ferromagnetism was observed in Co doped ZnO nanoneedle array prepared by pulsed laser deposition. L. Q. Liu et al.<sup>[7]</sup> reported the room temperature ferromagnetism of Zn<sub>1-x</sub>(FeCo)<sub>x</sub>O nanowires synthesized via a CVD growth method.

Recently, solution growth methods, for example, the bottom-up approach is becoming attractive to grow semiconducting magnetic nanostructures as they are expected to allow better control over chemical composition and dopant speciation compared to those of the physical methods. The chemical synthesis, in addition to being cost-effective, also offers advantages such as bulk synthesis and easy processing. Although there are a number of reports on the chemical synthesis of ZnO nanostructures, only few of them are related to the synthesis of doped nanostructure of ZnO. Yang et al.<sup>[8]</sup> reported the synthesis of Co-doped ZnO nanorods using the hydrothermal method. Wang et al. [9] reported a two-step chemical method for the synthesis of Co-doped singlecrystalline ZnO nanorods. The first step involved the preparation of ZnO nanoparticles (seeds), which acted as nucleating centers for nanorods, and in the second step nanorods were grown from the solution. In the chemical

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Corresponding author. Li Yingying, Ph. D., Associate Professor, Staff Room of Emergency Rescue, Public Security Fire Force College, Kunning 650208, P. R. China, Tel: 0086-871-67210444-8150, E-mail: waiwailyy@163.com

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methods reported so far, the final products were obtained in powder form. Hence, there is a need to further process these powders so as to prepare their films for making devices. Nowadays, substrate-anchored thin film with oriented nanostructures are often more desirable for applications. Three-dimensional (3D) arrays of nanostructured materials have been widely introduced to enable to improve unique performances in electronics <sup>[10]</sup>, optical devices <sup>[11]</sup>, light emitting diodes <sup>[12]</sup>, and so on. The low dimensional, wide band gap (3.37 eV) semiconductor-zinc oxide (ZnO) with large excition binding energy (60 meV) at room temperature has been so far applied in many aspects. Particularly, one dimensional ZnO oriented nanostructures are confirmed to be able to further enhance the performances of various devices, including field emission flat panel displays [13], shorewavelength lasers <sup>[14]</sup>, solar cell <sup>[15]</sup> and chemical sensors <sup>[16]</sup>. However, there are few researches on magnetic materials based on oriented ZnO nanostructure arrays. In particular, one-step template-free synthesizing route for one-dimensional ZnO array is still a significant challenge for researchers. In this paper, we reported a single-step simple wet chemical process to directly produce films of Co-doped ZnO nanorod arrays on zinc substrate in an aqueous medium. Zinc substrate can not only act as the carrier of the nanorods arrays, but also be the excellent conductor, which ensure the good ohmic contact between the ZnO nanorods and Zn substrate. At the same time, the ZnO nanorods film layer in-situ-generated on the zinc foil surface can improve the interface adhesive force.

#### 1 Experiment

All chemicals were analytical-grade reagents and purchased from Shanghai Chemical Reagent Corp. The zinc foil (99.99% purity) with a size of 2.5 cm×2.0 cm×0.1 cm was polished, washed by acetone, ethanol and distilled water for several times, and dried in an oven. Cobalt acetate [Co(Ac)<sub>2</sub> H<sub>2</sub>O] and tetramethylammonium hydroxide (TMAOH, 25%) were used as reactants without further purification.

The Zn foil substrates were put up-side down into a 50 mL Teflon-lined stainless steel autoclave within 40 mL aqueous solution of 0.3 mol/L TMAOH and required amount  $Co(Ac)_2 H_2O$ . The autoclave was heated to 170 °C for 12 h in the oven. After the reaction, the specimens were completely washed with absolute ethanol and Milli-Q water for several times, and then dried in vacuum at 60 °C for 8 h.

The morphologies of the as-synthesized Co-doped ZnO nanorods arrays were observed by field-emission scanning electron microscopy (FE-SEM, LEO1550). The crystal structures of the obtained samples were characterized by X-ray diffraction with graphite monochromatized CuK $\alpha$  radiation (BrukerD8,  $\lambda$ =0.154 05 nm). Photoluminescence measurement was carried out by employing a CARY Eclipse Fluorescence spectrum. The Raman spectra were collected at room temperature with a confocal Raman microscope

(NT-MDT NTEGRA Spectra) using a 514 nm laser as the excitation source. A vibrating sample magnetometer (VSM) was used to study the magnetic properties.

### 2 Results and Discussion

In the present experiment, we couldn't confirm the quantity of the zinc atoms involved in the reaction, so we assumed all the TMAOH reacted with Zn substrates and calculated the quantity of the  $Zn^{2+}$ . Then the Co(Ac)<sub>2</sub> H<sub>2</sub>O were added in the solution with the mass fractions 1%, 5%, and 10% of the Zn<sup>2+</sup>. The samples are denoted as ZnO:1wt%Co, and ZnO:5wt%Co, ZnO:10wt%Co.

Figs.1a~1g show the SEM images of the ZnO nanorods arrays with different Co doping concentrations and the pure ZnO nanorods arrays. Remarkably, the uniformity of the nanorods over a large area is obviously shown in a top view image, taken in a low magnification of the nanorods arrays (Fig.1a, 1c, 1e). Uniform growth of nanomaterials on large scale is important to nanodevice fabrication. The highmagnification FE-SEM images of the arrays (Fig.1b, 1d, 1f) suggest that the nanorods of high density grow nearly vertically from the substrates. The diameters of the nanorods in a hexagonal shape with the pencil shape tip are estimated to



Fig.1 SEM images of ZnO nanorods arrays with different Co doping concentrations and pure ZnO nanorods arrays:
(a, b) ZnO:1wt%Co; (c, d) ZnO:5wt%Co; (e, f) ZnO: 10%Co; (g) pure ZnO nanorods arrays

be 200~300 nm. The lengths of the nanorods are more than 5  $\mu$ m. Through comparing the Fig.1b, Fig.1d, Fig.1f and Fig.1g, it is found that the morphologies of the Co doped ZnO nanorods do not change significantly.

The XRD patterns of ZnO nanorods arrays doped with different concentrations of Co ions are shown in Fig.2. All diffraction peaks are well indexed to the standard diffraction pattern of hexagonal phase ZnO, implying a wurtzite structure with high crystallinity, without any additional impurity phases, thereby indicating that the Co ions successfully occupy the lattice sites rather than interstitial ones <sup>[17]</sup>. It also can be seen that with increasing of concentration of Co, the peak position of Zn<sub>1-x</sub>Co<sub>x</sub>O is shifted to higher angles, which indicates the decrease of lattice parameters *a* and *c* <sup>[18]</sup>. These phenomena presumably result from the substitution of Co ions with a small ionic radius of 0.058 nm for ZnO (0.060 nm) <sup>[19]</sup>. This also can be proved that the Co ions are replaced into the ZnO lattice.

Fig.3 presents the room temperature photoluminescence (PL) spectra of the as-grown Co-doped ZnO nanorods arrays and pure ZnO nanorods arrays by a Xe lamp (325 nm wavelength) as the source of excitation. The pure ZnO nanorods arrays exhibit a strong luminescence peak centered at 386 nm, which represents the near-band-edge emission. In addition, there is no other emission peak except a weaker blue luminescence emission peak around 443 nm. Compared with the undoped ZnO, the UV emission peak intensity decreases in Co-doped ZnO nanorods arrays with the increase of Co doping concentration, which perhaps are attributed to non-radiative recombination process. Besides, the band edge emission peak width FWHM (full width at half maximum) increases as the Co doping concentration increases. The broadening of the band edge PL emission may be explained in terms of potential fluctuations resulting from the random distribution of Co dopants in ZnO <sup>[20]</sup>. The microscopic inhomogeneity of the dopant concentration causes the potential fluctuation or tail states of band edges, leading to broadening of the band edge emission.



Fig.3 PL spectra of the as-grown Co-doped ZnO nanorods arrays and pure ZnO nanorods arrays, recorded at room temperature,  $\lambda_{ex}$ =325 nm

The pure ZnO nanorods arrays and Co-doped ZnO nanorods arrays have been investigated by Raman spectrum, as shown in Fig.4. For the pure ZnO, the sharpest and strongest peak located at 437 cm<sup>-1</sup> can be assigned to the high-frequency branch of  $E_{2H}$  mode of ZnO, which is the strongest mode in crystal structure. The peaks at 328 and 374 cm<sup>-1</sup> are assigned to the second-order vibration mode  $E_{2H}-E_{2L}$ of ZnO and the transverse-optical mode with A1 symmetry mode [A<sub>1</sub>(TO)], respectively. In contrast to Raman spectra of ZnO, the Raman-scattering peak positions of ZnO:1wt%Co, ZnO:5wt%Co and ZnO:10wt%Co do not change, which is consistent with XRD results, where Co doping does not change the samples structure <sup>[21]</sup>, and except for the character peak of ZnO, no other peaks related impurity phases are obtained. In addition, with the increment of doping content, the intensity of the peak at 570 cm<sup>-1</sup> increases, which is assigned to the longitudinal-optical mode with ZnO [A<sub>1</sub>(LO)]. Generally, this mode is induced by host lattice defects, such as oxygen vacancies and Zn interstitials. When Co<sup>2+</sup> is doped into ZnO lattice, oxygen vacancies and Zn interstitials are created. Furthermore, the concentration of these defects increases with



Fig.2 XRD patterns of the ZnO nanorods arrays with different Co doping concentrations



Fig.4 Raman spectra of the pure ZnO nanorods arrays and Co-doped ZnO nanorods arrays, recorded at room temperature,  $\lambda_{ex}$ =325 nm

increasing of the Co concentration, which is approved by the enhanced intensity of the peak at  $570 \text{ cm}^{-1}$  <sup>[22]</sup>.

The magnetic properties of the Co doped ZnO nanorods arrays were investigated using a vibrating sample magnetometer (VSM). Fig.5 shows the magnetic field dependence of the magnetization (M-H curves) for ZnO, ZnO:5wt%Co and ZnO:10wt%Co at room temperature. The diamagnetic contribution from the sample holder has been subtracted from the data. We can see that the pure ZnO nanorods arrays show diamagnetic behavior whereas all Co-doped ZnO nanorods show ferromagnetic behavior at room temperature. From the hysteresis loop, the saturation magnetization is found to increase with the increase of Co doping concentration. For ZnO: 10wt%Co sample, hysteresis loops with a coercive field  $H_{\rm c}$  of 52.8 kA/m can be observed at 300 K from curve c in Fig.5. The saturation magnetization of the sample is about  $6.105 \times 10^{-3}$  A m<sup>2</sup>/kg. There are two possible origins of the ferromagnetism in the Co doped ZnO system. One is due to the small secondary phases such as Co-metal clusters or Co-oxide precipitates. In the present studies, metallic Co and CoO are unlikely the source of this ferromagnetism because no detectable Co or CoO phases could be observed in the XRD patterns. According to the XRD patterns, absorption spectra and Raman spectra results, Co<sup>2+</sup> ions are successfully incorporated into the wurtzite lattice at the Zn<sup>2+</sup> sites, and then the possibility of Co precipitations can be ruled out. We suggest that the origin of the observed ferromagnetic order in the Co doped ZnO nanorods arrays can be explained using the carriers-mediated mechanism, by the RKKY type exchange mediated by the conduction carriers <sup>[23]</sup>. The conduction carriers in ZnO can be the native defects such as oxygen vacancy and interstitial zinc, which can be easily formed in doped ZnO. Recently, some works indicate that the as-grown ZnO nanowires are dominated by zinc interstitial, which is a shallow-level donor, and make the synthesized ZnO an n-type semiconductor <sup>[24]</sup>. As a result, the observed magnetic behavior may be attributed to the exchange interaction between free delocalized carriers (hole or electron from the valence band)



Fig.5 *M-H* curves of pure ZnO nanorods arrays and Co-doped ZnO nanorods arrays at 300 K

and the localized magnetic dipole moments of the magnetic ions (the localized d spins on the Co ions). The free carrier concentration has an effect on the ferromagnetism of  $Zn_{1-x}Co_xO$ . The Raman scattering study shows that defects in Co doped ZnO nanorods arrays increase with the increment of Co content (see Fig.4). Therefore, although more studies have to be done to understand the connection between the defects and the room-temperature ferromagnetism in transition metal-doped ZnO, present study has clearly shown that the defect is an important factor influencing the room temperature ferromagnetism.

#### 3 Conclusions

1) Large scale densely packed and vertically oriented pencil like nanorods arrays can be grown on zinc substrate.

2) The as-prepared products have wurtzite structure, and no other secondary phase is found in the nanorods.

3)  $Co^{2+}$  ions can substitute  $Zn^{2+}$  sites.

4) The doped nanorods arrays have obvious room temperature ferromagnetic behavior. The observed magnetic behavior may be attributed to the exchange interaction between free delocalized carriers (hole or electron from the valence band) and the localized magnetic dipole moments of the magnetic ions (the localized d spins on the Co ions).

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Li Yingying et al. / Rare Metal Materials and Engineering, 2016, 45(1): 0046-0050

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# 垂直生长于锌基底的 Co 掺杂 ZnO 纳米棒阵列显著室温铁磁性能

李莹滢<sup>1</sup>, 董 祥<sup>2</sup>, 张海黔<sup>3</sup>

(1. 公安消防部队昆明指挥学校,云南 昆明 650208)(2. 昆明理工大学 废弃物资源化国家工程研究中心,云南 昆明 650093)(3. 南京航空航天大学,江苏 南京 210016)

**摘 要:**通过简单的水热合成法在锌片基底上一步制备了 Co 掺杂的 ZnO 纳米棒阵列。纳米棒在基底上均匀分布,取向一致,垂直于基底大面积生长。样品结构均为六方纤锌矿结构,具有高结晶质量,不含其它杂相。随着 Co 掺杂浓度的增加,紫外发射峰强度逐渐下降,近带隙发射峰的半峰宽也较纯 ZnO 变宽。拉曼光谱显示 Co 的掺杂使纳米棒出现了氧空位和锌填隙本征缺陷。随着 Co 浓度的增加这些缺陷也随之增加。掺杂纳米棒阵列的磁滞回线表明样品具有明显的铁磁特征,并有较大的矫顽力 *H*<sub>c</sub>=52.8 kA/m。这种 ZnO 基稀磁半导体纳米棒阵列是一种在自旋电子器件中具有应用潜力的纳米材料。

关键词:半导体;磁性材料;纳米棒阵列;ZnO;水热法

作者简介: 李莹滢, 女, 1982 年生, 博士, 副教授, 公安消防部队昆明指挥学校训练部应急救援教研室, 云南 昆明 650208, 电话: 0871-67210444-8150, E-mail: waiwailyy@163.com